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INFLUENCE OF SCATTERING ON SEISMIC WAVES: VELOCITY AND ATTENUATION STRUCTURE OF THE UPPER CRUST IN SOUTHEAST MAINE

Edmund Reiter M. Nafi Toksöz Anton M. Dainty

Earth Resources Laboratory
Department of Earth, Atmospheric, and
Planetary Sciences
Massachusetts Institute of Technology
Cambridge, Massachusetts 02139

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INTRODUCTION

Short period Rayleigh waves (Rg) provide a convenient tool for studying the elastic properties of the upper crust. Group and phase velocities of Rg waves are particularly sensitive to variations in upper crustal shear wave velocity structure, and two station surface wave attenuation methods may be used to study the anelastic properties of the upper few kilometers of the crust. Despite many crustal geophysical studies, the shear wave velocity structure of the upper few kilometers of the crust remains relatively poorly constrained. Refraction experiments are generally concerned with deeper structure while reflection work usually involves only compressional velocities. Laboratory shear velocities using ultrasonic techniques have been measured for many of the rocks commonly found in the upper crust but may fail to adequately reproduce in situ conditions such as partial closure of cracks and incomplete saturation.

Near surface seismic sources such as quarry and refraction blasts are efficient generators of Rg waves. In New England, local networks archive hundreds of digital seismograms from quarry blasts every year. Typically the Rg phase is the largest amplitude arrival on the seismogram and is observable to approximately 100 km, depending on the size of the blast and the nature of the travel path. Recent crustal refraction surveys represent another sizeable source of digitally recorded short period Rayleigh waves. The density of receivers in a typical crustal refraction survey (one station per 1-2 km) provides an excellent geometry for correlating Rg velocities, and therefore shear velocity structure, with surficial sedimentary deposits and mapped bedrock units.

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A number of investigators have used Rg dispersion to study upper crustal velocity structure.

Kafka and Dollin (1985) used the dispersion of 0.5-2 Hz Rg waves to map lateral variations in

velocity structure in Connecticut. Båth (1975) observed normally dispersed Rg in the 0.5-2 Hz frequency range and concluded there was a general increase in velocity with depth in the upper 1-3 km of the crust. Anderson and Dorman (1973) studied the effects of surficial sedimentary deposits on 0.6-5 Hz Rayleigh waves in the New York City area. With the exception of Al-Husseini et al. (1981), previous studies involve sparse data sets with quarry blasts or small shallow earthquakes as sources and sparse local networks as receivers. Al-Husseini et al. (1981) calculated phase velocities of 10-30 Hz fundamental mode Rayleigh waves from f-k transforms of seismic reflection spreads in northern Saudi Arabia. The large volume of data allowed the mapping of 150,000 km² and correlation of phase velocity with various surficial sedimentary units.

This study focuses on determining the velocity and attenuation structure of the upper crust from Rg arrivals recorded in the 5-30 km distance range. Velocity profiles are obtained by an iterative, least squares, maximum likelihood inversion of group velocity observations; additional constraints on velocity structure include observation of Pg velocities. The density and geometry of group velocity observations in this study allow an in depth study of the behavior of velocity structure as a function of distance. Attenuation measurements using the two station method show relatively little scatter compared with other upper crustal attenuation studies (Pulli, 1983). Some problems with this method of determining attenuation with Rg waves are identified.

DATA

The data for this study are from a refraction survey conducted in Maine by the United States Geological Survey (Murphy and Luetgert, 1986,1987). The survey, conducted in the fall of 1984, was part of a broad-based geophysical investigation of the northern Appalachian Mountains that included seismic reflection, gravity and magnetic studies. Principal goals of these studies include an accurate mapping of Moho depth, identification of major intercrustal discontinuities and investigation of the complex region of a major plate collision that portions of Maine are commonly thought to represent.

The refraction survey consisted of 8 separate lines of 120 instruments, each approximately 90 km long and laid out as shown in Figure 1. Data from deployments 3, 4, 5 and 6 were obtained from the USGS and deployments 5 and 6 were selected for processing and study primarily on the basis of the authors' interests. The major structural feature of the study area is the Central Merrimack Synclinorium. Instrument and shot locations and elevations were determined using USGS 1:24,000 and 1:62,500 topographic maps and are assumed accurate to within 20 m. Shots consisted of 2000 lb of ammonium nitrate explosive located in a 20cm by 40m drill hole; shot times are assumed accurate to within \pm 2 msec. Figure 2 shows three seismograms representative of the data used in this study.

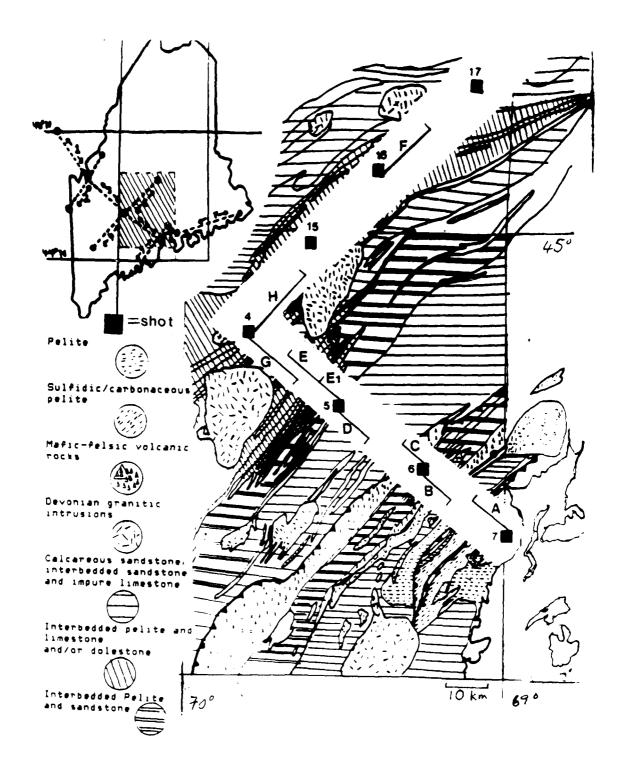


Figure 1. Map of general area and regions of study.

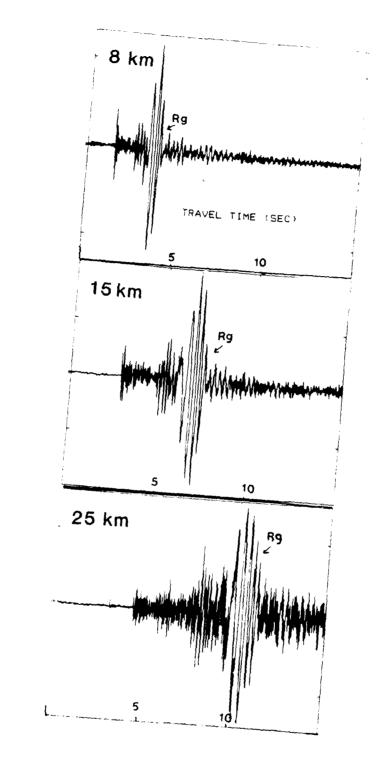


Figure 2. Typical seismograms recorded at 8, 15 and 25 km offsets. Amplitudes adjusted for display

ANALYSIS

Group velocities were determined using the multiple narrow bandpass filter technique described by Dziewonski and Hales (1972). Originally designed primarily for analysis of 20-200 second period Rayleigh waves, this technique worked well for the range of periods in this study. A brief description of the more important aspects of this technique is given below.

A Gaussian function

$$\mathcal{X}(\omega) = e^{-\alpha (\frac{\omega - \omega_c}{\omega_c})^2} \tag{1}$$

is chosen to represent the filter transfer function. This filter is chosen because of the mutual properties of its time and frequency durations. If we express the RMS time and frequency durations of a general signal as

$$(\mathcal{D}_t)^2 = \int_{-\infty}^{\infty} t^2 |f(t)|^2 dt \tag{2}$$

and

$$(\mathcal{D}_{\omega})^{2} = \int_{-\infty}^{\infty} \omega^{2} |F(\omega)|^{2} d\omega$$
 (3)

it can be shown that the value of the product of the RMS durations for the Gaussian function, $(\mathcal{D}_t \cdot \mathcal{D}_\omega)$, is the least of all the possible non-band limited functions, being equal to the constant $\sqrt{\pi/2}$ (Papoulis, 1962). This means that the frequency-time resolution, commonly measured as $1/(\mathcal{D}_t \cdot \mathcal{D}_\omega)$, is greater for a Gaussian filter than for any other non-band limited filter.

Representing the seismogram as

$$F(\omega) = A(\omega) \cdot e^{i(k(\omega)r - \omega t)} \tag{4}$$

and using first order Taylor expansions for $A(\omega)$ and $k(\omega)$ it can be shown that the maximum

amplitude of the filtered signal, given a filter of the form described above, corresponds to the group travel time of the energy associated with the frequency ω_c (Dziewonski et al., 1969). The seismogram is repeatedly filtered with appropriate values of ω_c and α and a group velocity dispersion curve derived from the peak amplitudes of the filtered envelopes.

Attenuation as a function of frequency was measured from the decay of Rg spectra recorded at distances between 10 and 30 km. The Rg arrival was windowed in the time domain and an amplitude spectrum calculated from the window. If $A(\omega, r_1)$ and $A(\omega, r_2)$ represent the amplitude spectra of two receivers at r_1 and r_2 , and $U(\omega)$ the interstation group velocity of the Rg wave, $Q(\omega)$ may be calculated as (Ben-Menahem and Singh, 1981)

$$Q(\omega) = \omega/(2 \cdot \gamma(\omega) \cdot U(\omega)) \tag{5}$$

with

$$\gamma(\omega) = \frac{\ln[A(\omega, r_1)/A(\omega, r_2)]}{(r_1 - r_2)} \tag{6}$$

Pg velocities were calculated by picking the arrival time of the phase, subtracting the known origin time to find the travel time and dividing the source-receiver distance by the travel time. Pg arrival times are accurate since at distances relevant to this study, 5-30 km, Pg is the first arrival and has an impulsive onset.

An iterative least squares maximum likelihood inverse routine was used to invert the group velocity observations for crustal velocity structure (Menke, 1984). A formulation by Rodi et al. (1975) was used to calculate partial derivatives of group velocity from partial derivatives of phase velocity. This formulation requires the calculation of a double set of roots and phase-velocity partials to obtain the group-velocity partials for all model parameters at a given frequency. It is

ultimately based on the derivative relationship between group and phase velocity. The necessary phase velocity, group velocity and phase velocity partial derivative calculations were carried out using techniques described in Takeuchi and Saito (1972). In this procedure variational principles are used to formulate expressions for group velocity and phase velocity partial derivatives which involve energy integrals and not numerical differentiation. This results in a more accurate determination of group velocity and a much quicker method of obtaining the phase velocity partial derivatives.

GROUP VELOCITY OBSERVATIONS

Approximately 400 seismograms were analyzed for Rg group velocity dispersion using the multiple filter technique discussed previously. As shown in Figure 1, source-receiver distances ranged from 5–30 km, and station spacing averaged one station per 800 m. We have chosen to focus on two features of the dispersion data. First, lateral variations in group velocity with group velocity differences of 0.25–0.35 km/sec are found and appear to correlate in some cases with mapped bedrock units and possibly tectonic features. Second, a dependence of group velocity on azimuth is observed. Pg velocities are consistent with both of the above observations. Before discussing these observations we will briefly examine some limitations and problems involved in the interpretation of Rg and Pg velocities.

Fundamental resolution problems involving primarily wavelength and receiver density limit the use of Rg waves in delineating regions of similiar dispersive properties. Examples presented here involve regions of roughly 10 km in length; assuming wavelengths ranging from 0.5-3 km, and receiver spacing of one recorder per 800 m, the wavelengths and receiver density are of proper size to adequately sample these areas. When interpreting Pg velocities for crustal velocity structure it is important to consider the difficulty in constraining the ray path of this phase. Pg ray paths in this distance range are very sensitive to the compressional velocity gradient in the upper 2 km of the crust, and given the lack of distinct refractor legs on the travel time curves calculation of absolute crustal velocities from these arrivals is difficult. However, by comparing Pg velocities calculated at similiar offsets in different regions the differences in ray paths is hopefully minimized, thus providing a reasonable quantity with which to compare velocity structures. In any event, if

the Pg velocities measured as described above are different at the same offset for two regions, we may certainly say that the structure is different.

Lateral Variation in Velocity Structure

Lateral variation in velocity structure was studied using group velocities measured for paths cutting the structural trend at approximately 90°. Group velocities for the six regions marked A, B, C, D, E and G were measured by calculating group velocities for three receivers located at the end of each boundary marked on Figure 1. Receiver spacing averaged about one receiver per 800 meters so the three group velocity measurements shown for each region are all located within a kilometer of the end of the marked boundaries.

Variations in Rg group velocity of 0.25-0.35 km/sec were observed between the fastest and slowest paths in this area (Figure 3). The fastest velocities were observed for region G which consists of about half calcareous sandstone and impure limestone, and half interbedded pelite and limestone. The slowest velocities were recorded about 50 km away in region B. This region is composed primarily of interbedded pelite and sandstone with minor amounts of mafic-felsic volcanics and calcareous sandstone. In regions G and B group velocities appear to associated with lithology but this is usually not the case. Regions A, D and E have nearly identical group velocities (Figure 3) yet have travel paths composed of distinctly different lithologies. Region A consists almost entirely of sulfidic/carbonaceous pelite while region D consists of calcareous sandstone and impure limestone. Region E is composed of about half interbedded pelite and sandstone and half interbedded pelite and dolestone. These three regions (A, D and E) all have group velocities ranging from around 2.55 km/sec at 1 Hz to 2.35 km/sec at 5 Hz yet include almost every lithology found

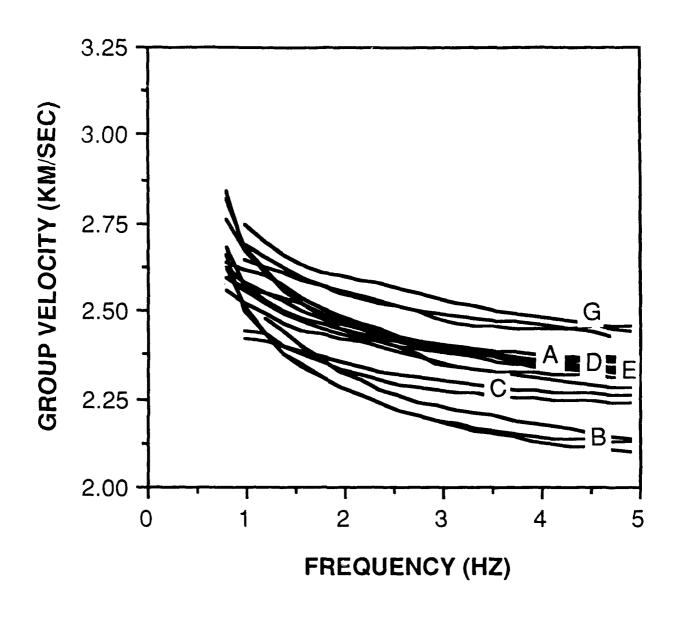


Figure 3. Group velocities for regions A, B, C, D, E and G. Note evidence for lateral variation in velocity structure.

in the region. Regions A and B are composed of different units and also are separated by a thrust fault. As is illustrated in Figure 3 these regions differ in group velocity by around 0.2 km/sec. This suggests that Rg group velocities might be helpful in delineating this tectonic boundary. In general, group velocities for cross strike paths vary from 2.4-2.6 km/sec at 1 Hz to 2.1-2.45 km/sec at 5 Hz. In some cases different lithologies are observed to display different Rg dispersive characteristics while in others the opposite is true.

Pg velocities were calculated for each of the regions discussed above. At the offsets relevant to these regions, 5-20km, Pg ray paths bottom at a maximum depth of around 2 km and especially at offsets near 20 km may be sampling slightly deeper velocity structure than the Rg waves. An average Pg velocity for a given region was determined by averaging Pg velocities determined at each of the three stations used in the previously discussed group velocity analysis. Table 1 lists these velocities and the offsets at which they were calculated. In general the Pg velocities show a similiar lateral variation in velocity structure but require further analysis such as ray tracing and travel time curve modeling in order to fully utilize this information. Such analysis is beyond the scope of this study and Pg velocities are presented only to provide a qualitative means of velocity structure comparison.

Azimuthal Dependence of Group Velocity

A dependence of Rg group velocity on azimuth was observed for all paths within the Central Merrimack Synclinorium. Faster velocities were observed for paths nearly parallel to the structural trend of the region and slower velocities for paths perpendicular to this direction. To illustrate this azimuthal dependence of group velocity, we will first examine four specific regions and then all

Table 1: Pg velocities showing differences between regions

Region	Offset (km)	Velocity (km/sec)
A	8	5.70
A	17	5.75
В	8	5.40
В	17	5.65
C	6	5.40
C	10	5.45
D	6	5.45
D	10	5.50
${f E}$	15	5.38
G	15	5.60

relevant data as a function of azimuth. Pg velocities will also be compared and used to substantiate the group velocity observations.

Figure 4 compares the group velocities for regions E and F. Both regions consist primarily of interbedded pelite, limestone and sandstone and have little (less than 150 m per 5 km) topographic variation. Thus, although over 90 km apart, regions E and F are of similiar composition and topography, and can be expected to exhibit similiar dispersion characteristics. Figure 4, however, shows a 15% difference in group velocities for the two regions. The group velocities measured for region E were calculated with an interstation method in order to more accurately measure the velocity of the pelite-limestone-sandstone unit by eliminating the effects of the portion of the path labeled E1. Pg velocities measured at a 15 km offset differed by 0.7 km/sec or 12%. However, the effect of E1 is present in the 15 km offset Pg velocity calculation. The group velocity of E1 is found to be slightly less than that of E, and thus is expected to bias the Pg velocity calculated in E downward, in effect increasing the difference between regions F and E. This effect is sufficiently

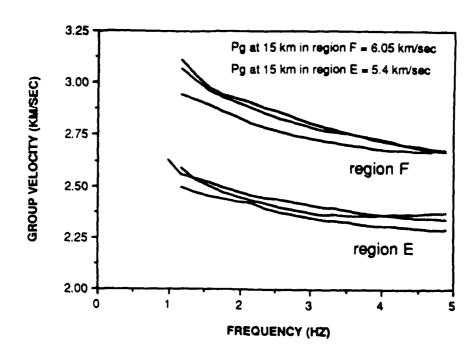


Figure 4. Group velocities for regions E and F.

small and well known to allow a valid comparison of Pg velocities between regions E and F and indicates that a difference of 10% is probably more representative of the actual in situ Pg velocities for these path geometries.

Group velocities for regions G and H are shown in Figure 5. As in all previous examples Pg velocities show the same trend as the Rg relocities; an Rg group velocity difference of 18% and a Pg velocity difference of 7% is observed for these two regions. These velocity differences between regions G and H provide some of the strongest evidence of an azimuthal dependence of velocity. As in the comparison of regions E and F, the mapped bedrock units in each region are quite similiar. However unlike regions E and F which are approximately 90 km apart, G and H are adjacent to one another. The proximity of G and H combined with the compositional and topographic similiarities suggest strongly that lateral variation in velocity structure is not the principal cause of the observed velocity differences.

Figure 6 is a plot of group velocity from several azimuths for paths entirely within the Central Merrimack Synclinorium. Although azimuthal coverage is sparse, group velocities measured from shot #16 in region F which cuts the structural grain at approximately 20° lie between the minimum (cross strike) and maximum (along strike) observations. Previous geophysical studies such as gravity surveys, seismic reflection and structural interpretations have failed to produce any evidence of the elongate, high velocity structure needed to explain the observed velocities with a laterally varying model. The apparent dependence of velocity on azimuth coupled with the lack of evidence for a laterally varying structure lead us to suggest a model requiring an anisotropic crust in at least the upper 2 km and probably significantly deeper. Symmetry axes of the anisotropy appear to

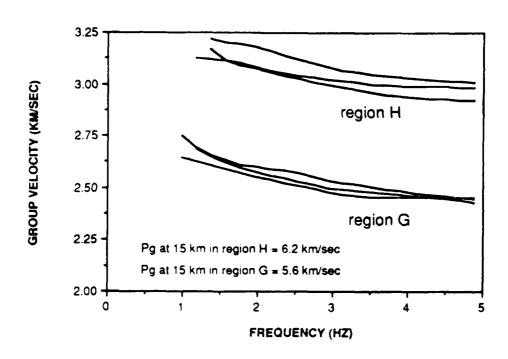


Figure 5. Group velocities for regions G and H.

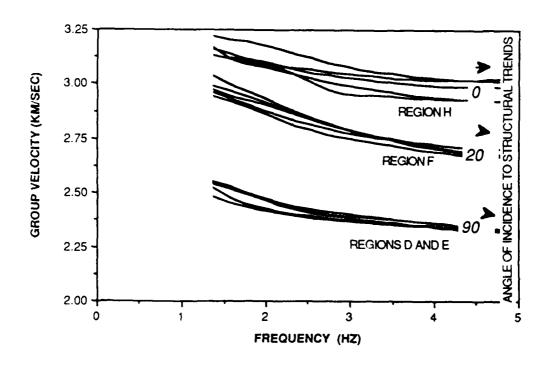


Figure 6: Group velocities from Central Merrimack Synclinorioum with angle of incidence with respect to gross structural trends

be determined by the orientation of the structural trend of the region which is approximately N30°E. Further analysis of the orientation of the symmetry axes involved is not possible with the azimuthally limited velocity observations of this study.

ATTENUATION OBSERVATIONS

The attenuation structure of the upper crust was studied using the surface wave $Q(\omega)$ analysis method discussed earlier. Difficulties with this method, arising primarily from suspected lateral variations/anisotropy in velocity structure, are identified and an average $Q(\omega)$ is determined from selected portions of the data set. $Q(\omega)$ calculated in this manner is found to be too unstable a quantity for use in delineating the individual attenuation properties of specific bedrock units, but provides a reasonable measure of the gross attenuation structure of the upper crust.

In the previously discussed formulation (5) and (6) for $Q(\omega)$ the implicit assumption is made that the difference between the two spectra $A(\omega, r_1)$ and $A(\omega, r_2)$, after correction for geometrical spreading, is due solely to the effect of dissipation of seismic energy. In a laterally homogeneous, isotropic, medium this is indeed the case and provided the exponential decay law is valid relations (5) and (6) provide a simple method for measuring $Q(\omega)$ from surface wave spectra. If, however, the structure has significant lateral velocity variation or anisotropy, the spectral difference observed between r_1 and r_2 is due not only to the effect of attenuation but also to the lateral change/anisotropy in velocity structure. Scattering from a sharp vertical velocity boundary (Chen and Alsop, 1979), fundamental to higher mode conversion (Kennett, 1984), multipathing (Yomogida 1985), and a different surface wave spectral response to the new medium must all be considered in the presence of lateral velocity variation and anisotropy. By averaging $Q(\omega)$ observations from many different regions we try to minimize the problem of lateral heterogeneity and feel reasonably confident that the $Q(\omega)$ values presented are representative of the gross attenuation properties of the upper crust, although the effect of anisotropy remains to be quantified.

An average $Q(\omega)$ was determined for regions J, K and L in Figure 7. Receiver separation averaged approximately 10 km with 9 station pairs used in the determination of an average Q for each region. The parts of the profiles labeled "1" in Figure 7 are the locations of the near receivers (r_2) ; the far receivers are located in the other part of the profiles. This source-receiver geometry has the effect of minimizing structural effects unique to an individual receiver while providing the redundancy necessary in attenuation studies. The spectra used for region J are shown in Figure 8. The investigation in this region suggests that lateral heterogeneities are a substantial problem when determining $Q(\omega)$ with the two station surface wave method. Group velocity curves for J and J1 are shown in Figure 8 and show a significant lateral variation in velocity structure. The dashed curve in Figure 8 is the group velocity curve used in the Q determination. It was determined by measuring the interstation group velocity between the three stations at approximately 15 km range and the three stations at approximately 25 km range. Also, there are large and very rapid variations of Q as a function of frequency for frequencies greater than 3.5 Hz. Such variations cannot be explained by a change of the average Q with depth because the Q at two nearly identical frequencies would be nearly the same, i.e., there is a limit to the rapidity of Q variations with frequency due to such a structure. For this reason we suspect that lateral hetrogeneity is playing a role.

Regions K and K1 (Figures 7, 9) represent another area in which lateral changes in velocity structure are suspected of causing substantial contamination of $Q(\omega)$ calculations. Group velocities measured for K and K1 again show a lateral variation in velocity structure. Also, spectra from the 14-16 km range (K1) lack the well developed "notch" near 2.5 Hz seen in spectra from the 24-26 km range. The large peak in Q at 2.5 Hz is caused by this difference, and for the same reasons as

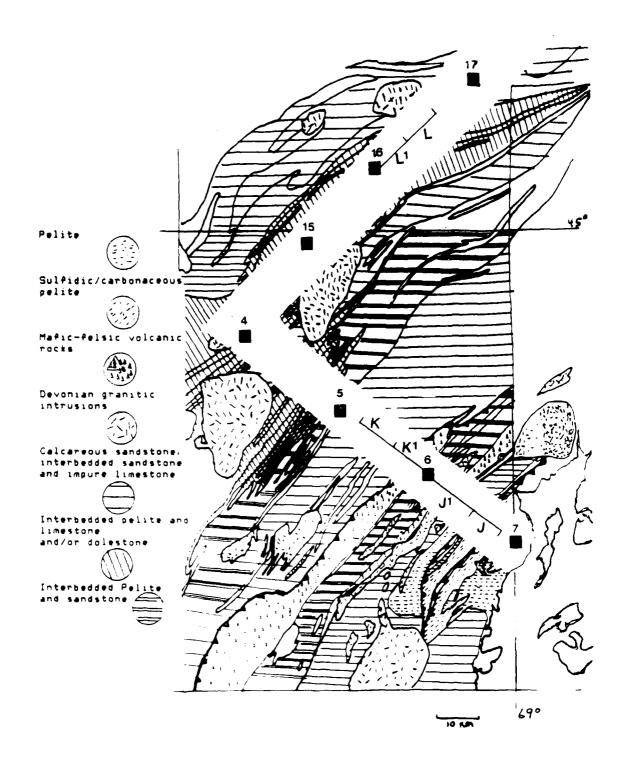


Figure 7. Map showing regions used in attenuation study and station geometry used in Q determinations.

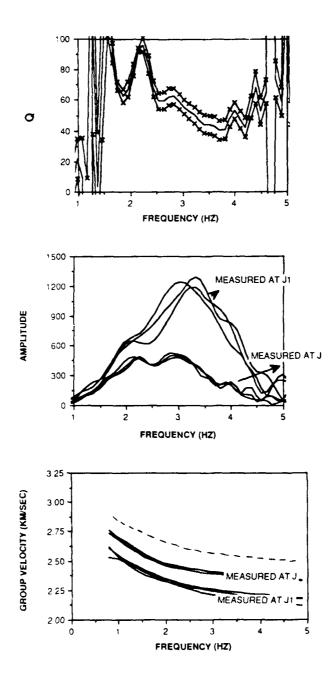


Figure 8. $Q(\omega)$, spectra and group velocities for regions J and J1. Dashed curve is the interstation group velocity used for Q determination.

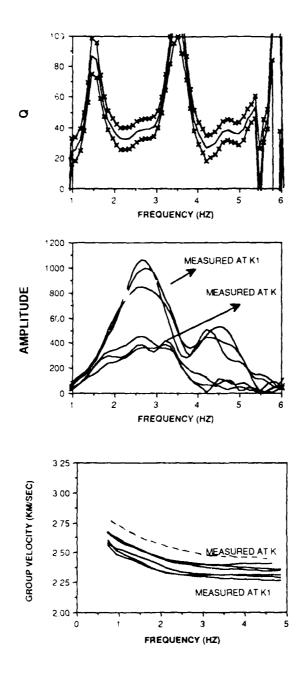


Figure 9. $Q(\omega)$, spectra and group velocities for regions K and K1. Dashed curve is the interstation group velocity used for Q determination.

before lateral heterogeneity is suspected.

An along strike attenuation curve was calculated for region L in Figure 7. In this $Q(\omega)$ determination an increase in Q is observed between 1 and 3 Hz and Q values of the same order as those observed for the cross strike regions are found (Figure 10). The increase in Q between 1 and 3 Hz could be due to lateral velocity variation, a frequency dependent $Q(\omega)$, or a low Q zone at depth. In theory a medium displaying an azimuthal dependence of velocity may also display an azimuthal dependence of attenuation. Although such behavior may exist, the effect is small enough to prevent unambigious detection in the presence of lateral variations in velocity structure.

Figure 11 shows an average $Q(\omega)$ from all regions except between shots #15 to #4 and #15 to #16 where the Rg waves were too poorly developed to measure either group velocity or Q. The overall attenuation structure of the upper crust leads to surface wave $Q(\omega)$ values between 80 and 25. Delineation of the Q properties of individual mapped bedrock units is not possible due to the effect of lateral variations in velocity structure which violate the approximation of homogeneity assumed in the surface wave Q method. A dependence of attenuation on azimuth is predicted to exist but its effect cannot be observed in the presence of the previously discussed difficulties in Q measurements.

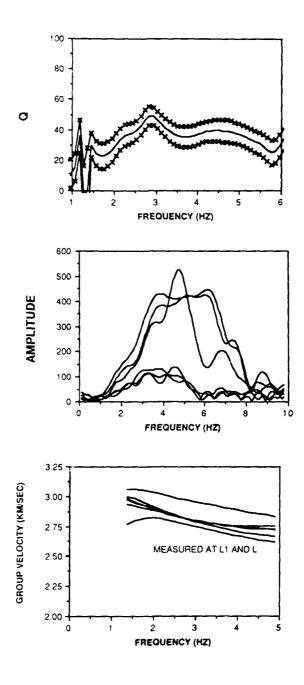


Figure 10. $Q(\omega)$, spectra and group velocities for regions L and L1.

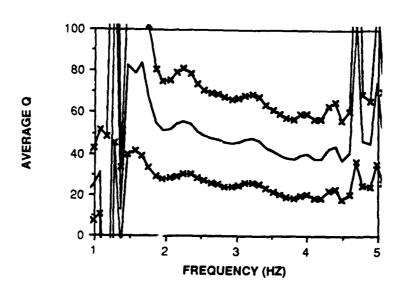


Figure 11. Average $Q(\omega)$ from all Q observations in this study.

INVERSION

A maximum likelihood inversion formulation discussed earlier is used to invert four sets of group velocity observations. This inversion assumes Gaussian distributions for the data observations and the initial model guess.

$$m_{k+1} = m_k + \left[A_k^T \cdot R_{dd} \cdot A_k + R_{mm}^{-1} \right]^{-1} \cdot \left[A_k^T \cdot R_{dd}^{-1} \cdot \left(d - G_k \cdot m_k \right) + \left(R_{mm} \right)^{-1} \cdot \left(m_0 - m_k \right) \right]$$
(7)

In (7), m_k is the kth model, A_k is the kth sensitivity matrix, R_{dd} , R_{mm} are the data and model envariance matrices, G_k is the forward model relationship, and d is the data vector.

A data vector consisting of twenty-two group velocity observations between 0.9 and 5 Hz and a model composed of 14 layers each containing two parameters, the compressional (V_P) and shear (V_S) velocities, describe the inverse problem. Layering varies from 25 m thick very near the surface to 200 m thick at 1 km to 500 m thick at 2 km. Group velocity observations from regions A and 3 (Figure 1), presented earlier as evidence of lateral variation in velocity structure, are inverted for shear and compressional velocity structure. An along strike velocity structure is obtained by inverting group velocity observations from region H, and a cross strike structure is obtained by inverting group velocity curves from region G (Figure 1). These inversions assume isotropy and thus may not accurately describe certain regions in this study.

Schematic partial derivatives of group velocity with respect to V_P and V_S , from the final velocity model for region A, are shown in Figure 12. The shear velocity derivative dominates at depths greater than about (wavelength/5) but the magnitude of the compressional velocity derivative is sufficient near the surface to warrant its inclusion in the formulation of the inverse problem. Figure 12 clearly illustrates the depth limitations on the velocity inversion imposed by the lowest

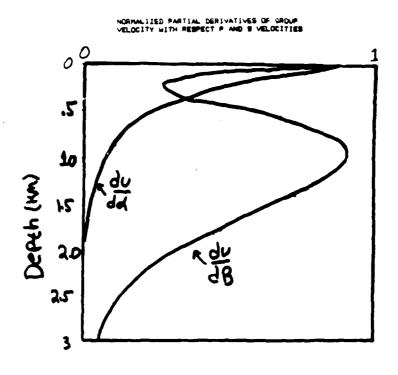


Figure 12. Partial derivatives of group velocity with respect to V_P and V_S for a 0.9 Hz Rayleigh wave using velocity model for region A.

observable frequency, in this case around 0.9 Hz.

The two group velocity curves from regions A and B and the inverted velocity models are shown in Figure 13. A shear velocity variation of 7 to 9% is found to exist between these two regions. Rows of the resolution matrix are plotted in Figure 14. A perfectly resolved parameter corresponds to a 1 on the diagonal and zeros elsewhere along the row corresponding to that parameter. The poorer the resolution the greater the deviation from this diagonalized delta function type behavior. The compressional velocities become poorly resolved at depth as is expected from Figure 12 and formally illustrated by Figure 14. The shear velocities are well resolved at all depths considered here. Compressional velocities determined from inversion of the group velocities agree well with in situ observations.

Along and cross strike velocity structures are found by inverting group velocities from regions G and H (Figures 15, 16). A difference of 15-20% is found to exist between these two velocity structures. A travel time curve was calculated from the compressional velocity model from the inversion and compared to the observed Pg travel time curve (Figure 17). The agreement between the predicted and the observed travel time curves is generally good. The theoretical travel time curve from the model for region G is systematically slower than the observations. The best fit compressional wave model for the data of Figure 17 is shown in Figure 15 as a dotted line and is compared to the compressional wave model from the surface wave inversion (solid line).

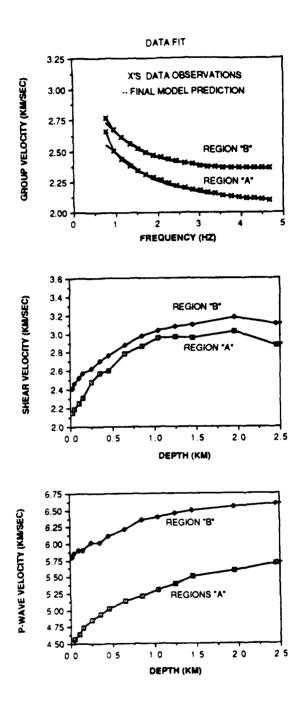


Figure 13. Velocity models for regions A and B and data fit. Each point in the velocity models corresponds to velocity in a given layer. Note evidence of lateral variation in velocity structure.

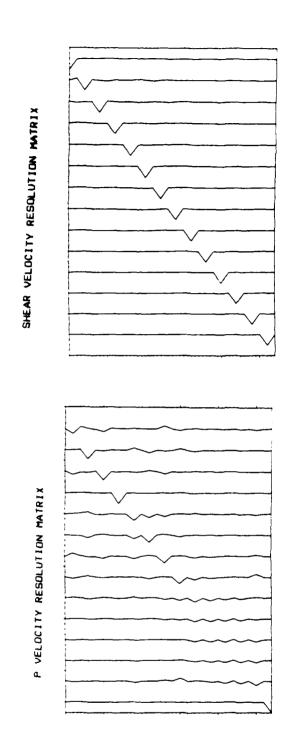


Figure 14. Resolution of shear and compressional velocity parameters, region A, represented as a plot of rows against columns of the resolution matrix. Each row corresponds to a model parameter, V_P or V_S , in a given layer. See Figure 13 for layer structure.

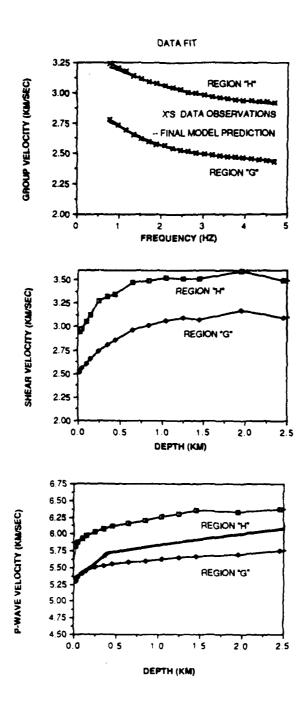


Figure 15. Velocity models for regions G and H and data fit. Dotted P-wave profile is derived by adjusting the velocity structure to match the observed Pg travel time curve. Note evidence of azimuthal variation of velocity. Each point in the velocity models corresponds to velocity in a given layer.

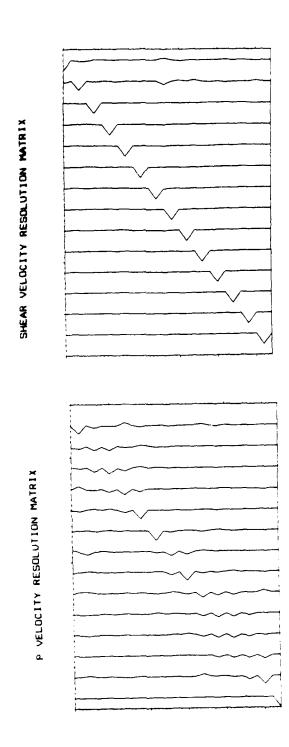


Figure 16. Resolution of shear and compressional velocity parameters, region H. Each row corresponds to a model parameter, V_P or V_S , in a given layer. See Figure 15 for layer structure.

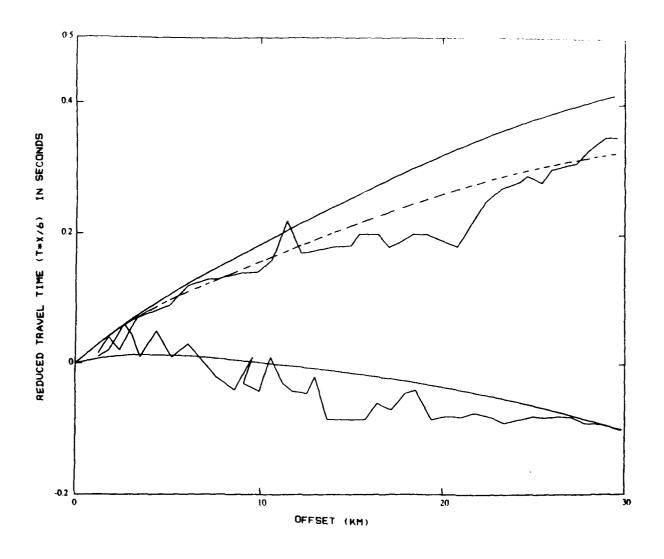


Figure 17. Comparison of observed Pg travel times with theoretical travel time curve from P-wave model from surface wave inversion (solid smooth curves). Dotted curve is best fitting model to travel time data after adjusting surface wave inversion model.

CONCLUSIONS

Lateral variations in velocity structure of up to 15% are observed from group velocity analysis of Rg waves and Pg velocities in southeastern Maine. Velocities determined from inversion of Rg group velocities agree generally with Pg velocities and the observed lateral variation in velocity correlates in some cases with mapped bedrock units and possibly tectonic features. A dependence of velocity on azimuth is observed for travel paths within the metasediments of the Central Merrimack Synclinorium. Insufficient azimuthal coverage prevents a more accurate description of this observation but the magnitude of the observed velocity difference, up to 20%, coupled with the geology of the area suggest that anisotropy in the metasediments is a likely source of this velocity behavior. Lastly, the attenuation structure of the upper two kilometers of the crust is characterized by Q values between 25 and 80. Lateral variations in velocity structure prevent precise determinations of Q when using two station surface wave Q methods, but given the difficulties involved in upper crustal attenuation studies Rg waves provide a relatively stable estimate of Q structure.

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- Yomogida, K. (1985). Amplitude and phase variations of surface waves in a laterally heterogeneous earth: ray and beam theoretical approach, *Ph.D. Thesis*, Massachusetts Institute of Technology

CONTRACTORS (United States)

Professor Keiiti Aki Center for Earth Sciences University of Southern California University Park Los Angeles, CA 90089-0741

Frofessor Charles B. Archambeau Cooperative Institute for Resch in Environmental Sciences University of Colorado Boulder, CO 80309

Dr. Thomas C. Bache Jr.
Science Applications I. '1 Corp.
10210 Campus Point Drive
San Diego, CA 92121 (2 copies)

Dr. Douglas R. Baumgardt Signal Analysis & Systems Div. ENSCO, Inc. 5400 Port Royal Road Springfield, VA 22151-2388

Dr. S. Bratt Science Applications Int'l Corp. 10210 Campus Point Drive San Diego, CA 92121

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Professor Robert W. Clayton Seismological Laboratory/Div. of Geological & Planetary Sciences California Institute of Technology Pasadena, CA 91125

Dr. Vernon F. Cormier
Department of Geology & Geophysics
7-45, Roon 207
The University of Conneticut
Storrs, Connecticut 06268

Dr. Zoltan A. Der ENSCO, Inc. 5400 Port Royal Road Springfield, VA 22151-2388

Professor John Ferguson Center for Lithospheric Studies The University of Texas at Dallas P.O. Box 830688 Richardson, TX 75083-0688 Professor Stanley Flatte'
Applied Sciences Building
University of California, Santa Cruz
Santa Cruz, CA 95064

Professor Steven Grand
Department of Geology
245 Natural History Building
1301 West Green Street
Urbana, IL 61801

Professor Roy Greenfield Geosciences Department 403 Daika Building The Pennsylvania State University University Park, PA 16802

Professor David G. Harkrider Seismological Laboratory Div of Geological & Planetary Sciences California Institute of Technology Pasadena, CA 91125

Professor Donald V. Helmberger Seismological Laboratory Div of Geological & Planetary Sciences California Institute of Technology Pasadena, CA 91125

Professor Eugene Herrin
Institute for the Study of Earth
& Man/Geophysical Laboratory
Southern Methodist University
Dallas, TX 75275

Professor Robert B. Herrmann
Department of Earth & Atmospheric
Sciences
Saint Louis University
Saint Louis, MO 63156

Professor Lane R. Johnson Seismographic Station University of California Berkeley, CA 94720

Professor Thomas H. Jordan Department of Earth, Atmospheric and Planetary Sciences Mass Institute of Technology Cambridge, MA 02139

Dr. Alan Kafka
Department of Geology &
Geophysics
Boston College
Chestnut Hill, MA 02167

Professor Leon Knopoff University of California Institute of Geophysics & Planetary Physics Los Angeles, CA 90024

Professor Charles A. Langston Geosciences Department 403 Deike Building The Pennsylvania State University University Park, PA 16802

Professor Thorne Lay
Department of Geological Sciences
1006 C.C. Little Building
University of Michigan
Ann Harbor, MI 48109-1063

Dr. Randolph Martin III New England Research, Inc. P.O. Box 857 Norwich, VT 05055

Dr. Gary McCartor Mission Research Corp. 735 State Street P.O. Drawer 719 Santa Barbara, CA 93102 (2 copies)

Professor Thomas V. McEvilly Seismographic Station University of California Berkeley, CA 94720

Dr. Keith L. McLaughlin S-CUBED, A Division of Maxwell Laboratory P.O. Box 1620 La Jolla, CA 92038-1620

Professor William Menke Lamont-Doherty Geological Observatory of Columbia University Palisades, NY 10964

Professor Brian J. Mitchell
Department of Earth & Atmospheric
Sciences
Saint Louis University
Saint Louis, MO 63156

Mr. Jack Murphy S-CUBED A Division of Maxwell Laboratory 11800 Sunrise Valley Drive Suite 1212 Reston, VA 22091 (2 copies) Professor J. A. Orcutt
Institute of Geophysics and Planetary
Physics, A-205
Scripps Institute of Oceanography
Univ. of California, San Diego
La Jolla, CA 92093

Professor Keith Priestley University of Nevada Mackay School of Mines Reno, NV 89557

Wilmer Rivers Teledyne Geotech 314 Montgomery Street Alexandria, VA 22314

Professor Charles G. Sammis Center for Earth Sciences University of Southern California University Park Los Angeles, CA 90089-0741

Dr. Jeffrey L. Stevens S-CUBED, A Division of Maxwell Laboratory P.O. Box 1620 La Jolla, CA 92038-1620

Professor Brian Stump Institute for the Study of Earth & Man Geophysical Laboratory Southern Methodist University Dallas, TX 75275

Professor Ta-liang Teng Center for Earth Sciences University of Southern California University Park Los Angeles, CA 90089-0741

Professor M. Nafi Toksoz
Earth Resources Lab
Dept of Earth, Atmospheric and
Planetary Sciences
Massachusetts Institute of Technology
42 Carleton Street
Cambridge, MA 02142

Professor Terry C. Wallace Department of Geosciences Building #11 University of Arizona Tucson, AZ 85721

Weidlinger Associates ATTN: Dr. Gregory Wojcik 620 Hansen Way, Suite 100 Palo Alto, CA 94304 Professor Francis T. Wu Department of Geological Sciences State University of new York At Binghamton Vestal, NY 13901

CTHERS (United States)

Dr. Monem Abdel-Gawad Rockwell Internat'l Science Center 1049 Camino Dos Rios Thousand Oaks, CA 91360

Professor Shelton S. Alexander Geosciences Department 403 Deike Building The Pennsylvania State University University Park, PA 16802

Dr. Ralph Archuleta
Department of Geological
Sciences
Univ. of California at
Santa Barbara
Santa Barbara, CA

Dr. Muawia Barazangi Geological Sciences Cornell University Ithaca, NY 14853

J. Barker
Department of Geological Sciences
State University of New York
at Binghamton
Vestal, NY 13901

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Dr. N. Biswas Geophysical Institute University of Alaska Fairbanks, AK 99701

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Professor Adam Dziewonski Hoffman Laboratory Harvard University 20 Oxford St. Cambridge, MA 02138

Professor John Ebel Dept of Geology & Geophysics Boston College Chestnut Hill, MA 02167

Dr. Alexander Florence SRI International 333 Ravenswood Avenue Menlo Park, CA 94025-3493

Dr. Donald Forsyth
Dept. of Geological Sciences
Brown University
Providence, Rl 02912

Dr. Anthony Gangi Texas A&M University Department of Geophysics College Station, TX 77843

Dr. Freeman Gilbert
Institute of Geophysics &
Planetary Physics
Univ. of California, San Diego
P.O. Box 109
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Mr. Edward Giller Pacific Seirra Research Corp. 1401 Wilson Boulevard Arlington, VA 22209 Dr. Jeffrey W. Given Sierra Geophysics 11255 Kirkland Way Kirkland, WA 98033

Dr. Henry L. Gray Associate Dean of Dedman College Department of Statistical Sciences Southern Methodist University Dallas, TX 75275

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Professor F.K. Lamb University of Illinois at Urbana-Champaign Department of Physics 1110 West Green Street Urbana, IL 61801

Dr. Arthur Lerner-Lam
Lamont-Doherty Geological Observatory
of Columbia University
Palisades, NY 10964

Dr. L. Timothy Long School of Geophysical Sciences Georgia Institute of Technology Atlanta, GA 30332

Dr. Peter Malin University of California at Santa Barbara Institute for Central Studies Santa Barbara, CA 93106

Dr. George R. Mellman Sierra Geophysics 11255 Kirkland Way Kirkland, WA 98033

Dr. Bernard Minster
Institute of Geophysics and Planetary
Physics, A-205
Scripps Institute of Oceanography
Univ. of California, San Diego
La Jolla, CA 92093

Dr. Geza Nagy SRI International 333 Ravenswood Avenue Menlo Park, CA 94025-3493

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Professor Larry J. Ruff Department of Geological Sciences 1006 C.C. Little Building University of Michigan Ann Arbor, MI 48109-1063

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1400 East 2nd South
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OTHERS (FOREIGN)

Dr. Peter Basham
Earth Physics Branch
Geological Survey of Canada
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Ottawa, Ontario
TANADA KIA 0Y3

Dr. Eduard Berg Institute of Geophysics University of Hawaii Honolulu, HI 96822

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Or. Kin-Yip Chun Geophysics Division Thysics Department University of Toronto Ontario, CANADA M5S 1A7

Ministry of Defense Blacknest, Brimpton, Reading RG7-4RS UNITED KINGDOM

Dr. Manfred Henger
Fed. Inst. For Geosciences & Nat'l Res.
Postfach 510153 D-3000 Hannover 51
PEDERAL REPUBLIC OF GERMANY

Dr. E. Husebye NTNF/NORSAR P.O. Box 51 N-2007 Kjeller, NORWAY Tormod Kvaerna NINE/NORSAR P.O. Box 51 N-2007 Kjeller, NORWAY

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Dr. Jack Evernden USGS-Earthquake Studies 345 Middlefield Road Menlo Park, CA 94025

U.S. Geological Survey ATTN: Dr. T. Hanks Nat'l Earthquake Resch Center 345 Middlefield Road Menlo Park, CA 94025

Dr. James Hannon Lawrence Livermore Nat'l Lab. P.O. Box 808 Livermore, CA 94550

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Paul Johnson ESS-4, Mail Stop J979 Los Alamos National Laboratory Los Alamos, NM 87545

Ms. Ann Kerr
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1400 Wilson Boulevard
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Dr. Kenneth H. Olsen Los Alamos Scientific Lab. Post Office Box 1663 Los Alamos, NM 87545

Howard J. Patton Lawrence Livermore National Laboratory P.O. Box 808, L-205 Livermore, CA 94550

AFOSR/NP
ATTN: Colonel Jerry J. Perrizo
Bldg 410
Bolling AFB, Wash D.C. 20332-6448

HQ AFTAC/TT Attn: Dr. Frank F. Pilotte Patrick AFB, Florida 32925-6001

Mr. Jack Rachlin - USGS - Geology, Rm 3 C136 Mail Stop 928 National Center Reston, VA 22092

Robert Reinke AFWL/NTESC Kirtland AFB, NM 87117-6008

HQ AFTAC/TGR Attn Dr. George H. Rothe Patrick AFB, Florida 32925-6001

Donald L. Springer Lawrence Livermore National Laboratory P.O. Box 808, L-205 Livermore, CA 94550

Dr. Lawrence Turnbull
OSWR/NED
Central Intelligence Agency
CIA, Room 5G48
Washington, D.C. 20505

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